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# On Design and Performance of Lossy Match GaAs MESFET Amplifiers

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**Abstract**—The noise figure, the gain, and the reflection coefficients of lossy match amplifiers and their dependence on circuit elements are studied. Theoretical results are supported by measured data taken on a two-stage unit. In addition, the performance characteristics of a 100-8800 MHz four-stage lossy match amplifier are discussed. The unit exhibits  $23.3 \pm 1.1$  dB of gain over the nearly 6-1/2 octaves. Its maximum noise figure is 10.9 dB from 100-8800 MHz and 6.6 dB from 2000-8000 MHz. The amplifier's overall circuit dimensions are  $10 \times 5.7$  mm.

## I. INTRODUCTION

THE BALANCED amplifier concept has dominated solid-state amplifier design for nearly two decades [1]. While it is expected to continue its dominating role, it will

not go unchallenged by three single-ended amplifier concepts that have made great progress in the field of broadband amplification over the last few years. These three principles, which are presently competing for the best single-ended amplifier performance, are [2]-[7] 1) the matched feedback amplifier; 2) the lossy match amplifier; and 3) the distributed amplifier. They all are characterized by their compact size and, at least the first two by their simplicity and low cost. It is for these reasons that the matched single-ended amplifier concepts are very attractive whenever an economical solution to wide-band amplification is of primary concern.

The application of dissipative gain compensation in interstage matching networks [8] and amplifier stabilization by means of resistive-loaded shunt networks [9] have been

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known for many years. Very recently analytic design techniques for an amplifier output network have been proposed [10]. It is the purpose of this paper to provide a better understanding of the electrical behavior of lossy match amplifiers and to study the influence of the individual circuit elements on the amplifier's performance characteristics such as noise figure, gain, and reflection coefficients. Formulas that express the important tradeoffs that exist between gain and reflection coefficients have been developed and are discussed. Furthermore, formulas for the amplifier's noise figure and minimum noise figure are presented. Finally, the experimental and theoretical results of a two-stage lossy match amplifier are compared and the measured performance of a compact ultrawide-band four-stage amplifier is reviewed.

## II. STUDY OF PERFORMANCE PARAMETERS

### A. Gain and Reflection Coefficients

The electrical behavior of an amplifier at microwave frequencies is best described by its scattering parameters which are a measure of the unit's insertion gain ( $|S_{21}|^2$ ), its input and output reflection coefficients ( $S_{11}$  and  $S_{22}$ ), as well as its reverse isolation ( $|S_{12}|^2$ ).

The schematic of a multistage lossy match amplifier as it is used in our experiments is shown in Fig. 1. This amplifier can be divided into three basic circuit functions: input matching, amplification, and interstage matching. The topology of the basic lossy match amplifier block containing dissipative circuit elements across its input and output terminals may be represented by the diagram shown in Fig. 2(a). The two-port between these two resistive shunt elements consists of the transistor and any additional elements that may be required for a desired performance. In the case of the basic amplifier block in Fig. 3(a), an open circuit shunt stub cascaded with a  $90\Omega$  transmission line at the transistor's drain terminal serves such a purpose.

Using the admittance representation of Fig. 2(a), one can easily derive the scattering parameters of this network. The results are presented as (A1) in the Appendix. Additional algebraic steps lead to an important relationship between  $S_{21}$  and the reflection coefficients  $S_{11}$  and  $S_{22}$  of the basic amplifier block

$$S_{21} = -\frac{Y_0}{Y_{12}} \left[ \sqrt{1 + \frac{Y_{21}Y_{12}}{Y_0^2} (1 + S_{11})(1 + S_{22})} - 1 \right]. \quad (1)$$

This expression leads to the calculation of the impact of any improvement of the reflection coefficients

$$S_{11} = \frac{[Y_0 - (Y_{11} + Y_G)][Y_0 + (Y_{22} + Y_D)] + Y_{12}Y_{21}}{[Y_0 + (Y_{11} + Y_G)][Y_0 + (Y_{22} + Y_D)] - Y_{12}Y_{21}} \quad (2a)$$

$$S_{22} = \frac{[Y_0 + (Y_{11} + Y_G)][Y_0 - (Y_{22} + Y_D)] + Y_{12}Y_{21}}{[Y_0 + (Y_{11} + Y_G)][Y_0 + (Y_{22} + Y_D)] - Y_{12}Y_{21}} \quad (2b)$$

( $Y_0$ —characteristic admittance)

by means of the admittances  $Y_G$  and  $Y_D$  on the unit's gain  $|S_{21}|^2$ .

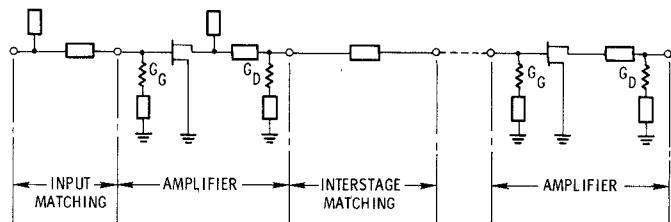


Fig. 1. Circuit topology of a multistage lossy match amplifier.

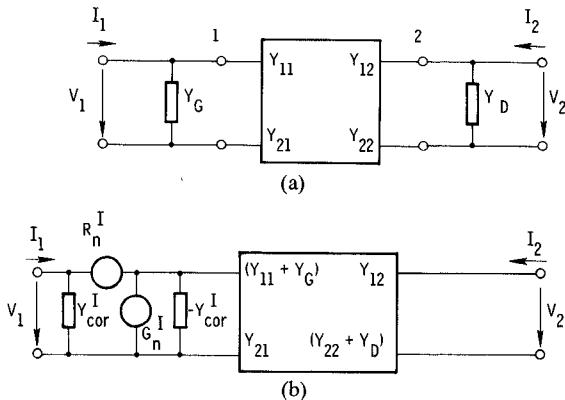
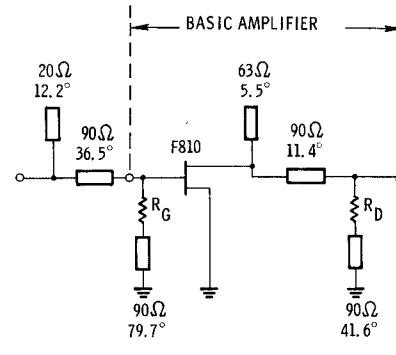


Fig. 2. Block diagram of the basic lossy match amplifier for the determination of (a) the scattering parameters and (b) the equivalent noise parameters.



(ELECTRICAL LENGTHS AT  $f = 8$  GHz)

(a)

### Equivalent Circuit Elements of WJ-F 810

$g_m = 57 \text{ mS}$	$R_g = 2.0 \text{ ohms}$
$\tau_0 = 4 \text{ p sec}$	$L_g = .208 \text{ nH}$
$C_{gs} = .508 \text{ pF}$	$R_s = 1.0 \text{ ohms}$
$C_{gd} = .041 \text{ pF}$	$L_s = .169 \text{ nH}$
$C_{dc} = .035 \text{ pF}$	$C_{ds} = .136 \text{ pF}$
$R_{gs} = 5.6 \text{ ohms}$	$R_d = 2 \text{ ohms}$
$R_{ds} = 257 \text{ ohms}$	$L_d = .146 \text{ nH}$

(b)

Fig. 3. Circuit topology of (a) a single-stage lossy match amplifier and (b) equivalent circuit elements of the GaAs MESFET.

In many practical cases the condition

$$\left| \frac{Y_{21}Y_{12}}{Y_0^2} (1 + S_{11})(1 + S_{22}) \right| \ll 1 \quad (3)$$

is satisfied and (1) adopts the much simpler form

$$S_{21} \approx -\frac{Y_{21}}{2Y_0}(1+S_{11})(1+S_{22}). \quad (4)$$

A close examination of (2) reveals that the input and output reflection coefficients  $S_{11}$  and  $S_{22}$  can only be made zero by means of the lossy shunt elements  $Y_G$  and  $Y_D$  as long as

$$G_{in} \leq Y_0 \quad (5a)$$

$$G_{out} \leq Y_0 \quad (5b)$$

where  $G_{in}$  and  $G_{out}$  are the conductances of the input and output admittances

$$Y_{in} = G_{in} + jB_{in} = Y_{11} + Y_G - \frac{Y_{12}Y_{21}}{Y_0 + Y_{22} + Y_D} \quad (6a)$$

$$Y_{out} = G_{out} + jB_{out} = Y_{22} + Y_D - \frac{Y_{12}Y_{21}}{Y_0 + Y_{11} + Y_G}. \quad (6b)$$

However, even if conditions (5) are not met, input and output matches may be improved by the reactive components of  $Y_G$  and  $Y_D$ .

In most practical cases of lossy match GaAs MESFET amplifiers, (5) is satisfied over most or at least over the lower portion of the frequency band. In order to improve the matches over the upper portion of the band, conventional matching techniques may be added by employing transmission line elements and open-circuit shunt stubs outside of the basic lossy match amplifier of Fig. 2(a). This may be accomplished in form of the input and interstage matching networks shown in Fig. 1.

The insertion gain of the basic amplifier derived from (2) is

$$\text{Gain} = \left| \frac{Y_0}{Y_{12}} \right|^2 \left| \sqrt{1 + \frac{Y_{21}Y_{12}}{Y_0^2} (1+S_{11})(1+S_{22})} - 1 \right|^2. \quad (7)$$

At frequencies where (3) is satisfied, the exact formula for gain (7) can be replaced by

$$\text{Gain} \approx \frac{1}{4} \left| \frac{Y_{21}}{Y_0} \right|^2 |1+S_{11}|^2 |1+S_{22}|^2. \quad (8)$$

The error when using (8) in case of the basic amplifier contained in Fig. 3(a) for  $R_G = R_D = 100 \Omega$  does not exceed 0.25 dB up to frequencies of 5 GHz.

While (7) expresses the tradeoffs between gain and reflection coefficients in its exact form, they can most easily be demonstrated on the basic amplifier's low-frequency model. For frequencies up to 1 GHz, the reactive elements of the transistor model listed in Fig. 3(b) have relatively little influence on the magnitude of the amplifier's gain and reflection coefficients. As a result, the transistor can be represented by its low-frequency model and the amplifier circuit can be reduced to the simple network inserted in Fig. 4. For this idealized model, the scattering parameters are easily determined. They are

$$S_{11} = \frac{1 - G_G Z_0}{1 + G_G Z_0} \quad (9a)$$

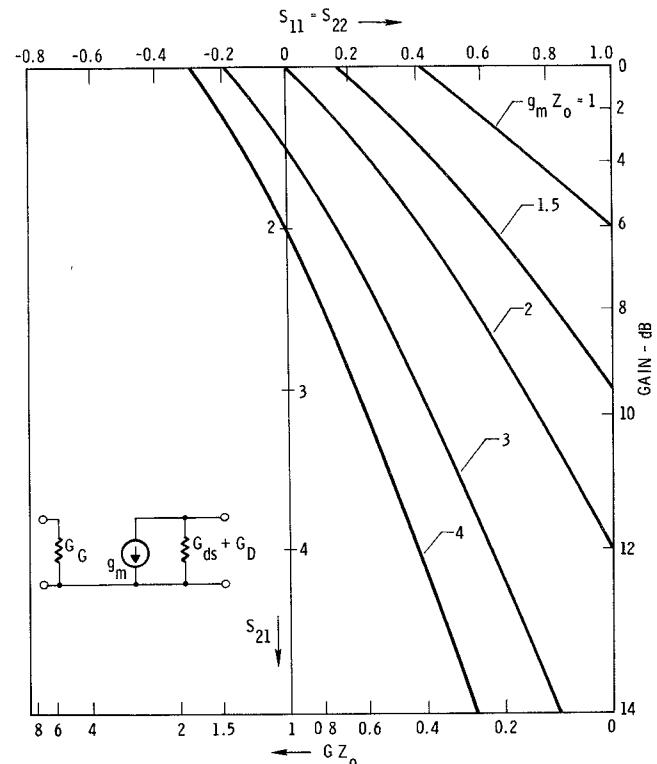


Fig. 4. Gain and reflection coefficients of a symmetric module ( $S_{11} = S_{22}$ ) at low frequencies as a function of  $GZ_0 = G_G Z_0 = (G_{ds} + G_D)Z_0$ .

and

$$S_{12} = 0 \quad (9b)$$

$$S_{21} = \frac{-2 g_m Z_0}{[1 + G_G Z_0][1 + (G_{ds} + G_D)Z_0]} \quad (9c)$$

$$S_{22} = \frac{1 - (G_{ds} + G_D)Z_0}{1 + (G_{ds} + G_D)Z_0} \quad (Z_0 - \text{characteristic impedance}). \quad (9d)$$

Furthermore, substituting  $Y_{21} = g_m$  and  $Y_{12} = 0$  into (1), where  $g_m$  is the transconductance, results in

$$S_{21} = -\frac{g_m Z_0}{2} (1 + S_{11})(1 + S_{22}). \quad (10)$$

At low frequencies,  $S_{11}$  and  $S_{22}$  have negligible imaginary components, and the gain can be expressed as

$$\text{Gain} = |S_{21}|^2 = \left[ \frac{g_m Z_0}{2} (1 + S_{11})(1 + S_{22}) \right]^2 \quad (11)$$

which clearly expresses the tradeoffs between the gain and the reflection coefficients. This relationship is presented in graphical form in Fig. 4 for a lossy match amplifier with identical input and output conductances; i.e.,  $G_G = G_{ds} + G_D = G$ , resulting in  $S_{11} = S_{22}$ .

### B. Noise Parameters and Noise Figure

The exact noise parameters of the basic lossy match amplifier [11] shown in Fig. 2(a) are

$$R_n^I = R_n + \frac{G_D}{|Y_{21}|^2} \quad (12a)$$

and

$$G_n^I = G_n + G_G + |Y_{11} - Y_{\text{cor}}|^2 \frac{G_D R_n}{|Y_{21}|^2 R_n + G_D} \quad (12b)$$

$$Y_{\text{cor}}^I = Y_{\text{cor}} + Y_G + (Y_{11} - Y_{\text{cor}}) \frac{G_D}{|Y_{21}|^2 R_n + G_D}. \quad (12c)$$

Their elements are located outside of the “noiseless” two-port as shown in Fig. 2(b).  $R_n$ ,  $G_n$ , and  $Y_{\text{cor}}$  are the equivalent noise resistance, equivalent noise conductance, and the correlation admittance, respectively of the two-port in Fig. 2(a) embedded between the shunt admittances  $Y_G$  and  $Y_D$ .

The exact minimum noise figure of our lossy match amplifier can now be calculated with

$$F_{\min}^I = 1 + 2 \left[ R_n^I G_{\text{cor}}^I + \sqrt{R_n^I G_n^I + (R_n^I G_{\text{cor}}^I)^2} \right]. \quad (13)$$

Fortunately, a number of conditions can be formulated that are satisfied in most cases of practical lossy match amplifier design and significantly simplify the resulting formula for the minimum noise figure. A condition met with few exceptions is

$$G_D \ll |Y_{21}|^2 R_n. \quad (14a)$$

If, in addition to (14a)

$$|Y_{11} - Y_{\text{cor}}|^2 G_D \ll |Y_{21}|^2 (G_n + G_G) \quad (14b)$$

and

$$G_{11} G_D \ll |Y_{21}|^2 R_n (G_{\text{cor}} + G_G) \quad (14c)$$

the minimum noise figure (13) may be expressed in the form

$$F_{\min}^I \simeq 1 + 2 R_n \left[ (G_{\text{cor}} + G_G) + \sqrt{\frac{(G_n + G_G)}{R_n} + (G_{\text{cor}} + G_G)^2} \right]. \quad (15)$$

Under the same conditions (14), the amplifier's noise figure [12], [13] is

$$F^I \simeq F_{\min}^I + \frac{R_n}{G_s^I} (G_s^I - G_{s\min}^I)^2 + \frac{R_n}{G_s^I} (B_s^I - B_{s\min}^I)^2 \quad (16)$$

with the source admittance

$$Y_s^I = G_s^I + jB_s^I \quad (16a)$$

and the source admittance for the minimum noise figure

$$Y_{s\min}^I = G_{s\min}^I + jB_{s\min}^I \simeq \sqrt{\frac{G_n + G_G}{R_n} + (G_{\text{cor}} + G_G)^2} - j \left( B_{\text{cor}} + B_G + B_{11} \frac{G_D}{|Y_{21}|^2 R_n} \right). \quad (16b)$$

### III. AMPLIFIER DESIGN

As is the case for all broad-band microwave solid-state amplifiers, the lossy match amplifier's performance improves with increasing transconductance and decreasing

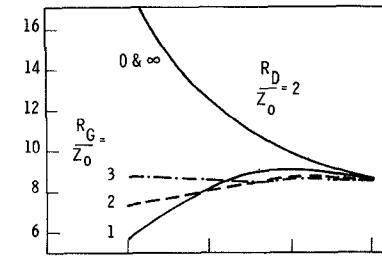


Fig. 5. The maximum available gain of the amplifier module (Fig. 3(a)) versus frequency with  $R_G/Z_0$  and  $R_D/Z_0$  as parameters.

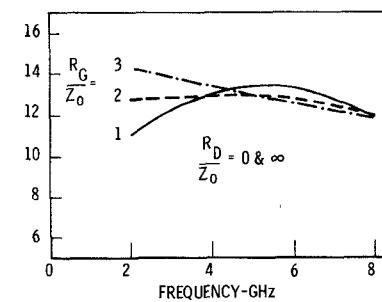


Fig. 6. The minimum noise figure of the amplifier module (Fig. 3(a)) versus frequency with  $R_G/Z_0$  and  $R_D/Z_0$  as parameters.

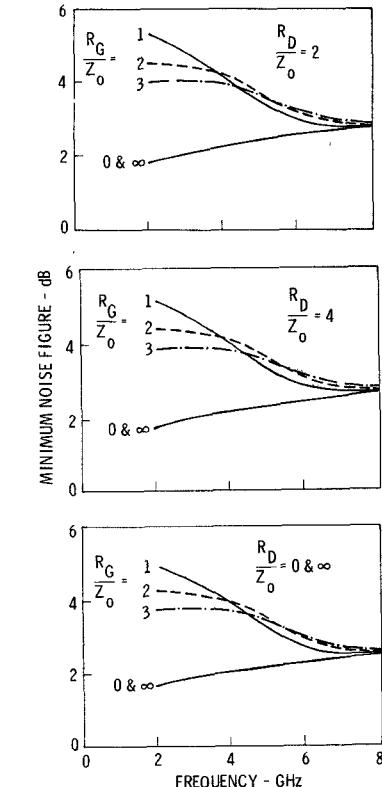


Fig. 6. The minimum noise figure of the amplifier module (Fig. 3(a)) versus frequency with  $R_G/Z_0$  and  $R_D/Z_0$  as parameters.

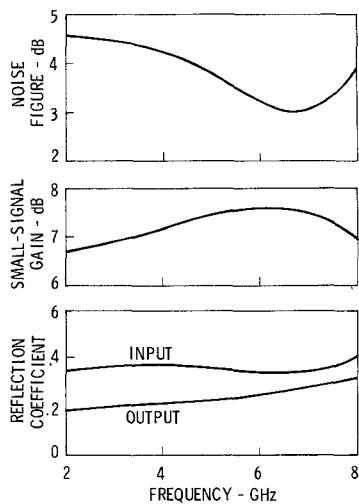


Fig. 7. Noise figure, gain, and reflection coefficients of the amplifier module (Fig. 3(a)) versus frequency with  $R_G = R_D = 100 \Omega$ .

parasitics of the GaAs MESFET. The normalized transconductance of the transistor as demonstrated by the curves in Fig. 4 should be  $g_m Z_0 = 2$  or greater in order to obtain reasonable gains and acceptable reflection coefficients. It is irrelevant whether high transconductance is achieved by one GaAs MESFET or several that are paralleled as long as parasitics don't limit the amplifier's performance at the high end of the desired frequency band.

The GaAs MESFET used in our experiments is the WJ-F810 whose equivalent circuit elements are listed in Fig. 3(b) and whose equivalent noise parameters have been published and discussed elsewhere [14]. The F810 has a normalized transconductance of  $g_m Z_0 = 2.85$ . Choosing  $G = G_G = G_{ds} + G_D = 10$  mS, one finds the gain and the reflection coefficients from Fig. 4:  $G = 8.1$  dB and  $S_{11} = S_{22} = 0.33$ . Taking into account that  $R_{ds} = 257$   $\Omega$ , the corresponding resistances are  $R_G = 100$   $\Omega$  and  $R_D = 164$   $\Omega$ . Since the drain current  $I_{ds}$  of the transistor flows through  $R_D$  it causes a significant voltage drop. In order to reduce this voltage drop and to simultaneously create better interstage matching conditions, it is advisable to reduce  $R_D$ . Using (9) and (11) one calculates  $G = 7.0$  dB,  $S_{11} = 0.33$ , and  $S_{22} = 0.18$  for  $R_G = R_D = 100$   $\Omega$ . Since gain decreases rapidly with frequency the resistive loss needs to be gradually reduced with frequency which is accomplished by the short-circuit shunt stubs in series with the resistors  $R_G$  and  $R_D$  of the circuit in Fig. 3(a). One finds that the electrical lengths of these high impedance lines is significantly less critical on the drain side of the amplifier. Introduction of the short-circuit shunt stubs significantly increases the gain over the second half of the frequency band and gain experiences a further improvement by placing the two circuit elements (see Fig. 3(a)) between the drain terminal and the resistive-loaded shunt circuit. Finally, the input reflection coefficient is reduced over the entire band by means of the simple input matching network shown in Fig. 3(a).

In order to demonstrate the influence of the resistors  $R_G$  and  $R_D$  on the module's optimum performance, the com-

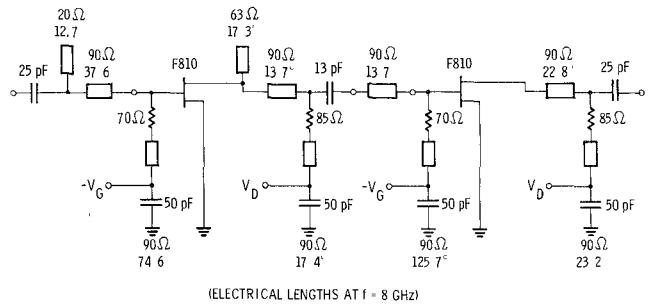


Fig. 8. Circuit topology of the computer optimized two-stage amplifier with the shunt resistors  $R_G = 70 \Omega$  and  $R_D = 85 \Omega$ .

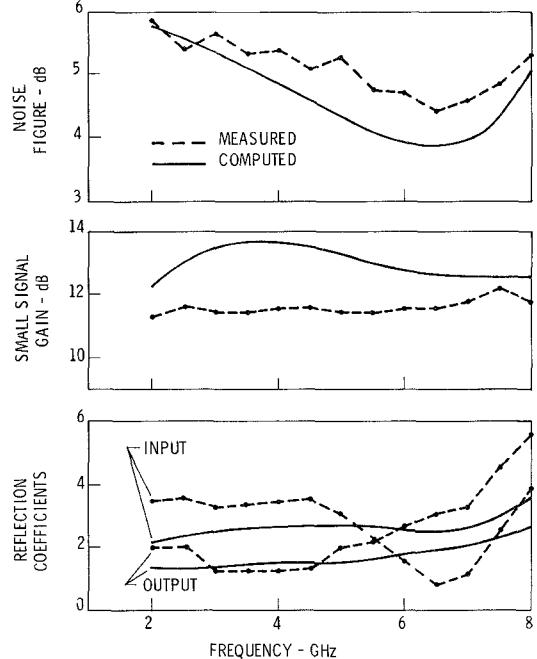


Fig. 9. Comparison of the computed and the measured performance characteristics of the two-stage amplifier of Fig. 8.

puted maximum available gain and minimum noise figure are plotted in Figs. 5 and 6, respectively. A comparison of the curves in Fig. 6 clearly illustrates that  $R_D$  has hardly any influence on the minimum noise figure of the amplifier module as one might expect for

$$\frac{G_D}{|Y_{21}|^2 R_n} \leq 0.22$$

across the 2–8-GHz band. The computed noise figure, gain, and reflection coefficients for  $R_G = R_D = 100 \Omega$  are plotted in Fig. 7.

#### IV. PERFORMANCE OF PRACTICAL AMPLIFIERS

When cascading lossy match amplifier modules, some of the unit's circuit elements require values somewhat different from those of the schematic in Fig. 3(a). This is due to the fact that the inputs and outputs of individual stages are no longer terminated by  $50\text{-}\Omega$  impedances.

The topology of a computer optimized two-stage amplifier is given in Fig. 8. The lossy match resistors  $R_G = 70 \Omega$  and  $R_D = 85 \Omega$  were not subjected to the optimization routine but rather determined experimentally by trying to

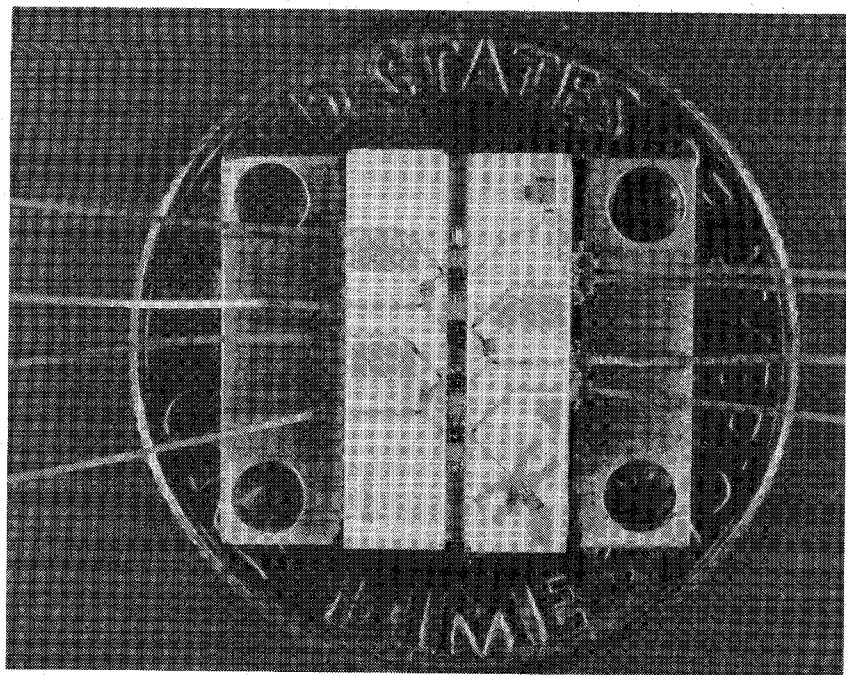


Fig. 10. Photograph of the four-stage amplifier.

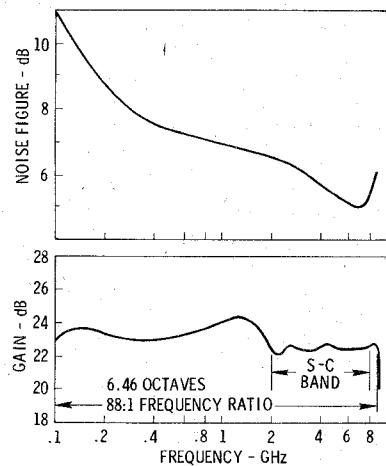
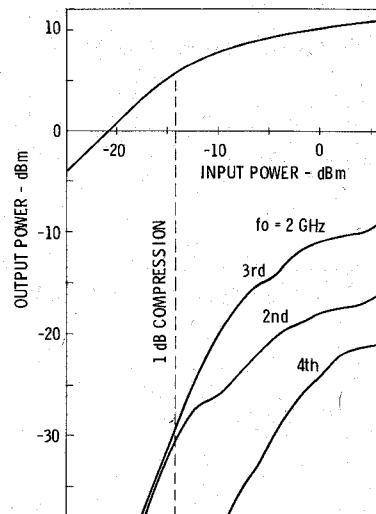


Fig. 11. Measured gain and noise figure of the four-stage amplifier.

achieve flat gain performance over the 2–8-GHz band. The measured and computed noise figures, gains, and reflection coefficients for this two-stage amplifier are compared in Fig. 9.

Increasing the number of stages from two to four and attempting to achieve flat gain performance required an additional reduction in the lossy match resistances to  $R_G = 60 \Omega$  and  $R_D = 80 \Omega$ . In addition, an increase in the impedances of the amplifier's transmission line elements from  $90 \Omega$  to  $115 \Omega$  became necessary. A photograph of this amplifier is shown in Fig. 10. Realized on 0.025 in. alumina, the overall circuit length of the unit is 0.394 in. or 10 mm. The unique layout of the circuit limits the number of substrates to two regardless of the number of stages and allows the transistors to be die attached to a single cooling rib. Measured gain and noise figures of this amplifier are plotted in Fig. 11. The results demonstrate the enormous bandwidth capacity of the lossy match principle. Extending

Fig. 12. Fundamental and harmonic output power of a four-stage amplifier ( $V_{DS} = 4$  V,  $I_{DS} = 35$  mA each).

the band to still lower frequencies is merely a question of using higher value bias and blocking capacitors. The output versus input curves of fundamental and harmonic output power of an earlier four-stage unit for  $f_0 = 2$  GHz are plotted in Fig. 12 when operating each transistor at 4 V of drain-source voltage and 35 mA of drain-source current.

## V. CONCLUSION

The noise figure, the gain, and the reflection coefficients of the lossy match amplifier have been studied and formulas to calculate these parameters were developed. Based on the theoretical results a two-stage amplifier was designed and tested. The unit's computed performance was compared with actual test data over the design frequency band of 2–8 GHz. The comparison demonstrates good

agreement between the measured and computed results.

Finally, the measured performance characteristics of a four-stage ultrawide-band amplifier were discussed. This unit has a gain of  $G = 23.3 \pm 1.1$  dB and a maximum noise figure of  $F = 10.9$  dB over the frequency band of 100–8800 MHz representing nearly 6-1/2 octaves of instantaneous bandwidth. Realized on only two substrates, the overall circuit dimensions are  $10.0 \times 5.7$  mm with 2.3 dB/mm of gain per unit length. These results evidence the enormous bandwidth potential of the lossy match amplifier achievable in an extremely compact size.

## APPENDIX

The basic circuitry of a lossy match amplifier can be represented by a two-port with lossy elements across its input and output ports as shown in Fig. 2. The two-port between ports 1 and 2 contains the transistor and additional lossless elements that may be required for a desired performance.

The circuit is then characterized by the admittance matrix

$$\begin{bmatrix} I_1 \\ I_2 \end{bmatrix} = \begin{bmatrix} Y_{11} + Y_G & Y_{12} \\ Y_{21} & Y_{22} + Y_D \end{bmatrix} \begin{bmatrix} V_1 \\ V_2 \end{bmatrix} \quad (A1)$$

which can easily be converted into the scattering matrix  $S_{ij}$ , whose elements are

$$S_{11} = \frac{[Y_0 - (Y_{11} + Y_G)][Y_0 + (Y_{22} + Y_D)] + Y_{12}Y_{21}}{[Y_0 + (Y_{11} + Y_G)][Y_0 + (Y_{22} + Y_D)] - Y_{12}Y_{21}} \quad (A2a)$$

$$S_{12} = \frac{-2Y_{12}Y_0}{[Y_0 + (Y_{11} + Y_G)][Y_0 + (Y_{22} + Y_D)] - Y_{12}Y_{21}} \quad (A2b)$$

$$S_{21} = \frac{-2Y_{21}Y_0}{[Y_0 + (Y_{11} + Y_G)][Y_0 + (Y_{22} + Y_D)] - Y_{12}Y_{21}} \quad (A2c)$$

$$S_{22} = \frac{[Y_0 + (Y_{11} + Y_G)][Y_0 - (Y_{22} + Y_D)] + Y_{12}Y_{21}}{[Y_0 + (Y_{11} + Y_G)][Y_0 + (Y_{22} + Y_D)] - Y_{12}Y_{21}} \quad (Y_0 - \text{characteristic admittance}) \quad (A2d)$$

Using the  $S$ -parameters of (A2) one obtains

$$S_{21} = \frac{-Y_{21}}{Y_0 + Y_{11} + Y_G} (1 + S_{22}) = \frac{-Y_{21}}{Y_0 + Y_{22} + Y_D} (1 + S_{11}) \quad (A3)$$

suggesting a tradeoff that exists between reflection coefficients and insertion gain. This tradeoff becomes more evident when substituting for  $(Y_{11} + Y_G)$  and  $(Y_{22} + Y_D)$  obtaining

$$S_{21} = -\frac{Y_{21}}{2Y_0} \frac{(1 + S_{11})(1 + S_{22})}{1 - \frac{S_{12}}{2} \frac{Y_{21}}{Y_0}} \quad (A4)$$

and with

$$S_{21}Y_{12} = S_{12}Y_{21}. \quad (A5)$$

(A4) can finally be expressed as a function of  $Y_{21}$  and  $Y_{12}$  and the reflection coefficients of the amplifier  $S_{11}$  and  $S_{22}$

$$S_{21} = -\frac{Y_0}{Y_{12}} \left[ \sqrt{1 + \frac{Y_{21}Y_{12}}{Y_0^2} (1 + S_{11})(1 + S_{22})} - 1 \right]. \quad (A6)$$

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# Weakly Coupled Dielectric Resonators

JEAN VAN BLADEL, FELLOW, IEEE

**Abstract** — The resonant modes of a pair of coupled resonators of high  $\epsilon_r$  are considered in the limit of large spacings  $D$  between resonators. Attention is focused on the lowest "magnetic-moment" mode, where the coupling effect leads to a split of the original mode into an even and an odd part. Formulas are obtained for the coupling coefficient, the resonant frequencies and  $Q$  of the modes. They are strikingly similar to those for weakly-coupled  $R-L-C$  circuits. The accuracy of the formulas is verified by comparing their predictions with direct numerical data, available for coupled circular cylindrical resonators.

## I. INTRODUCTION

THE RESONANT modes of a dielectric resonator of high  $\epsilon_r$  are obtained by solving the differential problem

$$\begin{aligned} -\operatorname{curl} \operatorname{curl} \bar{h}_m + k_m^2 \bar{h}_m &= 0 && \text{in the dielectric} \\ \operatorname{curl} \bar{h}_m &= 0 && \text{outside the dielectric.} \end{aligned} \quad (1)$$

The eigenvector  $\bar{h}_m$  must be continuous across the boundary surface, and of order  $1/R^2$  (or higher) at large distances [1]. The numerical solution of (1), a difficult three-dimensional problem, simplifies when the body is of revolution (e.g., a sphere or a circular cylinder [2]). In the coupled structure shown in Fig. 1, (1) must be solved in the presence of a composite dielectric consisting of volumes 1 and 2. When the spacing between 1 and 2 is small, the field distribution in each volume will be significantly perturbed with respect to that of the resonator isolated in free space. Any simplifying feature, such as symmetry of revolution, will be lost, except in structures of the type shown in Fig. 4.

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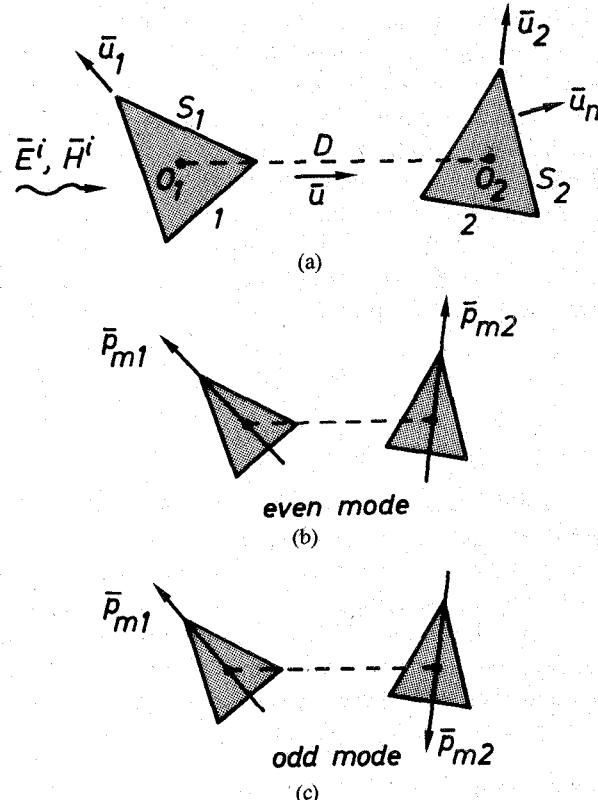


Fig. 1. Coupled resonators with corresponding resonant modes.

The conclusion is clear: the three-dimensional problem must be solved. To avoid this very arduous task, approximations have been made in the past [3], [4]. Cohn, for example, represents dielectric resonators by conducting loops carrying currents  $I$ , and endowed with  $L$ ,  $C$ , and